QTAIM-DI-VISAB computational study on the so-called nonclassical bicyclobutonium cation

Nick H. Werstiuk,^a* Daniel A. Poulsen^b

^aDepartment of Chemistry, McMaster University Hamilton, ON L8S 4M1, Canada ^bDepartment of Chemistry, University of California Berkeley, CA 94720-1460 *E-mail: werstiuk@mcmaster.ca*

This paper is dedicated to Ted Sorensen for his contributions to chemistry on the occasion of his 75th birthday

Abstract

QTAIM-DI-VISAB analyses at the CCSD and B3PW91 levels were used to characterize the bonding of the cyclopropylcarbinyl (1) and the so-called 'nonclassical' bicyclobutonium and 1-methylbicyclobutonium cations, 2 and 3 as well as the transition state for rearrangement of 1 to 2. These analyses involved obtaining QTAIM molecular graphs and delocalization indexes (DIs) for pairs of atoms that were correlated with the proximities of atomic basins (VISAB). This study showed that the supposed nonclassical bicyclobutonium and 1-methylbicyclobutonium cations do not exhibit penta-coordinate carbons at their equilibrium geometries as has been claimed. Both species are best described as distorted cyclobutyl cations that exhibit a single ring critical point in the topology of the charge density.

Keywords: Cyclopropylcarbinyl cation, bicyclobutonium cation, 1-methybicyclobutonium cation, DFT, MP2 and CCSD calculations, QTAIM-DI-VISAB analysis, molecular structure

Introduction

The cyclopropylcarbinyl cation (1) and its isomer, the so-called nonclassical bicyclobutonium cation (2) as well as its 1-methyl analogue (3) whose structures are shown using the usual dashed-line formalism have been the focus of numerous experimental and computational studies over the span of several decades with the latest being the work of Olah and coworkers.¹ The dashed-line structures of 2 and 3 have gradually been replaced with ORTEP-type/solid-line structures, shown as 4, the implication being that C3 is a penta-coordinate atom in both representations.



Numerous computational studies have been carried out over the space of two decades ranging from MINDO/3² to MP4(SDTQ)/MP2//MP2/6-31G(d,p) + ZPE³ and MP4(SDTQ)/cc-pVTZ//MP2/cc-pVTZ + ZPE levels¹. Computational studies were also carried out in solvent fields.⁴ The main focus of the most recent study on **2** and **3** carried out by Olah et al. at various levels, including MP4(SDTQ)/cc-pVTZ//MP2/cc-pVTZ + ZPE was the calculation of ¹³C NMR chemical shifts at CCSD(T).¹ With the exception of Brown's proposal⁵ based on solvolysis data the cation should be viewed as a distorted cyclobutyl cation represented as **5**, dashed-line or solid-line structures have been used to represent the bonding of these alleged nonclassical pentacoordinate species in all the studies carried out to this point.

Our recent computational studies on a number of cations, including 2-norbornyl, established that coordination based on the number of bond paths – as defined in a OTAIM 6 molecular graph - terminating at a nucleus in any species - cation, carbanion, radical, or carbene - should be used as the criterion of hyper-coordination and hyper-valency.^{7,8,9,10} A QTAIM treatment precludes interpretational bias of the nature of bonding based on qualitative evidence such as interatomic distances. We argued that this approach should be used regardless of the nature of the intermediate to obviate the confusion and inaccuracies associated with arbitrarily using indicators such as dashed lines, dotted lines, cross-hatched lines, hollow tubes, and solid tubes in structural formulas. In addition to using QTAIM molecular graphs that unambiguously define molecular structure, we recently showed that QTAIM-DI-VISAB analyses are useful for characterizing the bonding in molecules at their equilibrium geometries and refining our knowledge of the bonding in hyper-coordinated species.^{11,12,13,14} In addition to obtaining molecular graphs, delocalization indexes (DIs) are used to establish the importance of delocalization between pairs of atoms that are correlated with the proximity of atomic basins as visualized in displays of these basins (VISAB). This paper reports the results of a QTAIM-DI-VISAB study on 1, 2, and 3, the goal being to establish whether 2 and 3 are pentacoordinate species as has been claimed.

Computational methods

Our previous experiences with DFT calculations on carbocations clearly showed that the B3PW91 hybrid functional is superior to B3LYP in computing the geometries of delocalized, socalled nonclassical species.^{7,8,9,10} We obtained additional support for this finding by carrying out calculations on O-protonated 2,2-dimethyloxirane – studied by Carlier et al.¹⁵ and described as a particularly challenging computational problem - to compare results from B3LYP, B3PW91, PBE1PBE, and CCSD calculations at the 6-311G(d,p) level as implemented in G03.¹⁶ The results presented in an earlier publication¹³ clearly established that B3PW91 and PBE1PBE are expected to be superior to B3LYP for obtaining equilibrium geometries in cases where relatively weak polar bonds are involved. In this study, cation geometries were optimized at B3PW91 and MP2 levels with a range of basis sets - 6-311+G(2d,p), cc-pVTZ, and aug-cc-pVTZ with OPT=TIGHT and OPT=VERYTIGHT and wave functions obtained. The OPT=VERYTIGHT B3PW91 calculations on 2 were carried out with INT=ULTRAFINE. On the other hand, CCSD(full) calculations with OPT=TIGHT were carried out first with 6-31+G(d,p) and then with the 6-311+G(2d,p) basis set. It should be pointed out that while MP2(frozen core) geometry optimizations with the cc-pVTZ – as reported by Olah et al.¹ – and aug-cc-pVTZ basis sets converged, this was not the case at MP2(full); we found that none of the cc-pVTZ and aug-ccpVTZ MP2(full) calculations converged. Consequently, the MP2 calculations were discounted and not used in this study. Moreover, it appears that MP2 often tends to overestimate the stability of 'nonclassical' structures.¹⁷ While 1 and 2 were viewed as unsymmetrical species by Casanova et al.⁴, we found that OPT=TIGHT calculations yielded geometries that were extremely close to the C_{S} structures. Consequently we fixed the symmetry – in our case using Chemcraft¹⁸ – to C_{S} as did Olah et al.¹

Selected inter-nuclear distances of $1-C_S$, $2-C_S$, and $3-C_S$ obtained at CCSD(full)/6-311+G(2d,p) with OPT=TIGHT are displayed as Figure 1(a), (c), and (d), respectively. Frequency calculations were carried out on the stationary points to confirm them as energy minima or transitions states. B3PW91 and CCSD energies and thermochemical data are collected in Table 1. QTAIM analyses of the wave functions to investigate the topologies of the electron densities were carried out with AIM 2000¹⁹ and the obtained molecular graphs are shown in Figure 2 to 4. AIMALL 08²⁰ was used to integrate atomic basins, obtain atomic populations, calculate total charges as well as atomic overlap matrices required for DI calculations. Values of $\rho(\mathbf{r}_{c})$ at selected bond critical points are collected in Table 2. That the total charges of the cations obtained at the various levels of theory differed by less than 1% from the expected value of 1.0 confirmed the quality and validity of the QTAIM data (Table 3). This conclusion was supported by the fact that $E_{elec}(\Sigma E(\Omega))$, the sum of the atom energies $E(\Omega)$ obtained with AIMALL for 1, **TS-1** \rightarrow **2**, **2**, and **3** differed from the G03 molecular energy E_{elec} (Table 3) by less than 0.45 kcal mol⁻¹. The program LI-DICALC^{21,22,23} was used to obtain DIs; selected values for pairs of atoms are collected in Table 4. Isosurface plots of the density of atomic basins (Figure 5 -Figure 8) were obtained with AIM 2000 at a contour value of 0.005 which includes > 95% of the electrons using a mesh grid size of 0.125 and plotted with a sphere size of 0.15.

Results and Discussion

Thermochemistry

Table 1. Total and relative energies of cations at B3PW91/aug-cc-pVTZ(tight) and CCSD(full)/6-311+G(2d,p)(tight)

Cation	1- <i>C</i> _S	TS-1→2	$2-C_S$	3- <i>C</i> _{<i>S</i>}
$E_{elec}{}^{a}$	-156.290401	-156.289613(+0.48) ^h	-156.291993(<i>1.00</i>) ^h	-195.624835
		$(-136.8 \text{ cm}^{-1})^{i}$		
$E_{elec}{}^{\rm b}$			-155.291993	
$E_{elec}^{c}(CCSD)$	-155.973458	-	-155.973170(+0.18)	-195.227514
E_0^{d}	-156.193468	-156.192121(+1.09)	-156.193856(-0.24)	-195.499150
E ₂₉₈ ^e	-156.188554	-156.188239	-156.189409	-195.493515
$H_{298}{}^{f}$	-156.187609	-156.187296 (+1.96)	-156.188460(-0.54)	-195.492571
$G_{298}{}^{g}$	-156.220706	-156.217907(+1.76)	-156.220120(+0.37)	-195.527173

^a E_{elec} is the uncorrected total energy in hartrees. ^b E_{elec} is the uncorrected total energy in hartrees at B3PW91/aug-cc-pVTZ(opt=verytight, int=ultrafine). ^c E_{elec} in bold is the uncorrected CCSD(full)/6-311+G(2d,p)(tight) total energy in hartrees. ^d $E_0 = E_{elec} + ZPE^{e}E = E_0 + E_{vib} + E_{rot} + E_{trans}$. ^f H = E + RT. ^g G = H + S. ^h Values in brackets relative to 1- C_S in kcal mol⁻¹. ⁱ The imaginary frequency.

Selected thermochemical data for 1-*C_S*, TS-1-2, 2-*C_S*, and 3-*C_S* are collected in Table 1. As was found previously¹ cations 1-*C_S*, and 2-*C_S* are very close in energy: at B3PW91/aug-cc-pVTZ, 2-*C_S* is lower in energy than 1-*C_S* based on the ZPE-corrected difference in energy ΔE_0 (-0.24 kcal mol⁻¹) and ΔH_{298} (-0.54 kcal mol⁻¹). On the other hand 1-*C_S* was found to be lower in energy than 2-*C_S* based on ΔG_{298} (+0.37 kcal mol⁻¹) at B3PW91/aug-cc-pVTZ(tight) and ΔE_{elec} (+0.18 kcal mol⁻¹) at CCSD(full)/6-311+G(2d,p)(tight). As mentioned in the computational methods section we chose to discount MP2 calculations because the optimizations did not converge at the MP2(full)/cc-pVTZ or MP2(full)/aug-cc-pVTZ levels. Given the expected prohibitive length of CCSD(full)/6-311+G(2d,p) transition state and frequency calculations we studied TS-1-2 only at the B3PW91/aug-cc-pVTZ(tight) level. As seen in the data collected in Table 1, the vales for $\Delta E_{elec}^{\dagger}$, ΔE_{0}^{\dagger} , ΔH_{298}^{\dagger} , and ΔG_{298}^{\dagger} for the rearrangement of 1-*C_S* are +0.48, +1.09, +1.96, and +1.76 kcal mol⁻¹, respectively, consistent with the fact that values ranging from 0.60 to 2.2 kcal mol⁻¹ – uncorrected and ZPE corrected – reported by Olah et al. at various high levels of theory.¹

Equilibrium geometrical and molecular structures

Selected equilibrium internuclear distances of $1-C_S$, TS-1 $\rightarrow 2$, $2-C_S$, and $3-C_S$ are collected in Figure 1.



Figure 1. (a) selected internuclear distances (Å) of $1-C_s$ at CCSD(full)/6-311+G(2d,p)(tight), (b) TS-1 \rightarrow 2 at B3PW91/aug-cc-pVTZ(tight), (c) $2-C_s$ at CCSD(full)/6-311+G(2d,p)(tight), and (d) $3-C_s$ at CCSD(full)/6 311+G(2d,p)(tight).

The data obtained at B3PW91/aug-cc-pVTZ differed somewhat from the CCSD(full)/6-311+G(2d,p) values. For example the crucial C1-C3 distances of **2** and **3** are 1.651 and 1.745 Å, respectively, at B3PW91/aug-cc-pVTZ compared to 1.675 and 1.746 Å as shown in Figure 1, suggesting that like MP2 – Olah et al. reported values of 1.648 and 1.698 Å for the C1-C3 distances of **2**- C_S and **3**- C_S at MP2(fc)/cc-pVTZ¹ – B3PW91 overestimates the delocalization between C1 and C3 of **2**- C_S and **3**- C_S somewhat. Even at B3PW91/aug-cc-pVTZ with opt=verytight and int=ultrafine the distances the C1-C3 distance of **2**- C_S is 1.651 Å. Olah et al. reported values of 1.717 and 2.446 Å for the C1-C4 and C2-C4 distances of **TS-1→2** at MP2(fc)/cc-pVTZ¹ compared to 1.706 Å and 2.245 Å as shown in Figure 1.

Cation			<u>_p(</u>	<u>r) BCP</u>				<u>ρ(r)</u> <u>RCP</u>
	C1-C2	C2-C3	C1-C3	C1-C4	C2-C4	C3-C4	C1-C5	
1- <i>C</i> _S	0.3435 ^a	0.1656	-	-	_ ^d	0.2959	-	0.1631
	(0.3515) ^b	(0.1759)				(0.3515)		(0.1737) ^b
TS-1→2	(0.3407)	(0.2402)	-	NO	NO	(0.2709)	-	NO RCP
				BCP	BCP			
	0.2878	0.1753	NO BCP	_ ^d	-	_d	-	0.1571
$2-C_S$	(0.2954)	(0.1815)	NO BCP					(0.1618)
	(0.2954) ^c	(0.1815)	NO BCP					(0.1618)
3- <i>C</i> _{<i>S</i>}	0.2841	0.1843	NO BCP	_ ^d	-	_ ^d	0.2627	0.1338
	(0.2901)	(0.1929)	NO BCP				(0.2736)	(0.1379)

Table 2. Selected QTAIM data for cations at CCSD(full)/6-311+G(2d,p)(tight) andB3PW91/aug-cc-pVTZ(tight)

^a At CCSD(full)/6-311+G(2d,p)(tight).^b Values in italics in brackets at B3PW91/aug-ccpVTZ(tight). ^c Values in bold italics in brackets at B3PW91/aug-cc-pVTZ(opt=verytight, int=ultrafine).^d Data for atom pairs related by symmetry not included.

Of importance is the fact that the C1-C3 distance of **3** is 0.071 Å greater than the C1-C3 distance of **2** at CCSD(full)/6-311+G(2d,p). Selected QTAIM data for cations at CCSD(full)/6-311+G(2d,p)(tight) and B3PW91/aug-cc-pVTZ(tight) are collected in Table 2. The molecular graph of **1** obtained at CCSD(full)/6-311+G(2d,p)(tight) is displayed as Figure 2(a). The Poincaré-Hopf relationship (NumNACP + NumNNACP - NumBCP + NumRCP - NumCCP = 1) is satisfied. The C2-C3 and C2-C4 bond paths (BPs) are highly curved and the BCPs are in close proximity to the RCP; this means that rearrangement to another molecular structure by coalescence of a BCP and the RCP, as for example through **TS-1→2**, would require small nuclear displacements and a small activation energy in keeping with the very low barrier calculated for the rearrangement of **1** to **2**. Its molecular structure does not exhibit BCPs between C1|C4 nor C2|C4, and no RCPs.

Table 3. QTAIM charges and energies at CCSD(full)/6-311+G(2d,p)(tight) and B3PW91/aug-cc-pVTZ(tight)

Cation	Charge	$E_{elec}(\Sigma E(\Omega))$	E _{elec}
1- <i>C</i> _S	1.0003 ^a	-155.973727	-155.973458
	(0.9997) ^b	(-156.290405)	(-156.290401)
TS-1→2	(1.0001)	(-156.289480)	(-156.289613)
	1.0021	-155.973847	-155.973170
$2-C_S$	(1.0012)	(-156.292354)	(-156.291993)
	(1.0015) ^c	(-156.292467) ^a	(-156.291993)
$3-C_S$	1.0005	-195.227542	-195.227514
	(1.0008)	(-195.625037)	(-195.624835)

^a At CCSD(full)/6-311+G(2d,p)(tight). ^b Values in italics in brakets at B3PW91/aug-cc-pVTZ(tight). ^c Values in bold italics in brackets at B3PW91/aug-cc-pVTZ(opt=verytight, int=ultrafine).



Figure 2. (a) molecular graph of $1-C_S$ at CCSD(full)/6-311+G(2d,p)(tight), (b) **TS-1→2** at B3PW91/aug-cc-pVTZ(tight). Black spheres are carbon atoms and grey spheres are hydrogens. The red spheres are (3,-1) bond critical points (BCPs) and the yellow sphere is the (3,+1) ring critical point (RCP).

The molecular graph of 2- C_S obtained at CCSD(full)/6-311+G(2d,p)(tight) is displayed as Figure 3(a) and 3(b) with 3(b) clearly showing the puckered nature of the ring structure. The Poincaré-Hopf relationship is also satisfied in the case of 2- C_S . The C2-C3 and C3-C4 bond paths (BPs) exhibit considerable curvature. The most important point is that 2- C_S does not exhibit a BCP/BP between C1 and C3 at its equilibrium geometry at the CCSD(full)/6-311+G(2d,p) level; C3 is not a penta-coordinate carbon atom – it does not have five bond paths terminating at the nucleus. This is the case at the B3PW91/aug-cc-pVTZ(tight) level and B3PW91/aug-cc-pVTZ(opt=verytight, int=ultrafine) as well.



Figure 3. (a) molecular graph of $2-C_s$ at CCSD(full)/6-311+G(2d,p)(tight) and (b) $2-C_s$ at CCSD(full)/6-311+G(2d,p)(tight) re-orientated.

The molecular graph of $3-C_s$ obtained at CCSD(full)/6-311+G(2d,p) is displayed as Figure 4(a) with 4(b) showing the puckered nature of its ring structure. The Poincaré-Hopf relationship is satisfied. The C2-C3 and C3-C4 bond paths (BPs) don't exhibit the same degree of curvature

as the corresponding BPs of $2-C_S$ indicating qualitatively that $3-C_S$ has a more stable molecular structure than $2-C_S$. Once again the important point is that, like $2-C_S$, $3-C_S$ does not exhibit a BCP/BP between C1 and C3 at its equilibrium geometry at the CCSD(full)/6-311+G(2d,p) level; *C3 is not a penta-coordinate carbon atom*. Like $2-C_S$, this is also the case at the B3PW91/aug-cc-pVTZ(tight) level.



Figure 4. (a) molecular graph of **3**- C_S at CCSD(full)/6-311+G(2d,p)(tight), and (b) **3**- C_S at CCSD(full)/6-311+G(2d,p)(tight) re-orientated.

QTAIM-DI-VISAB Analyses

1-*C*_S

Selected atomic basins of 1-*C_s* obtained at the CCSD(full)/6-311+G(2d,p)(tight) level are displayed as Figure 5(a) and (b); Figure 5(a) shows that the C1 and C3 basins do not significantly impinge on each other in accord with the fact that they that exhibit a miniscule DI of 0.087: there is little delocalization of electrons between the C1 and C3 basins. As expected the DI (0.141) is somewhat larger at B3PW91/aug-cc-pVTZ. The C1-C2 bond exhibits considerable double bond character – $\rho(\mathbf{r})_{cp}$ is 0.3435 and the DI 1.317. On the other hand the C2-C3 and C2-C4 bonds are very weak – $\rho(\mathbf{r})_{cp}$ is 0.1656 and DI 0.590 at CCSD(full)/6-311+G(2d,p). The relative areas of the atomic surfaces shared with C2 seen in Figure 5(a) are in accord with the relative strengths of the C1-C2 and C2-C3 bonds. These values are proportionally larger (see Table 2 and Table 3) at B3PW91/aug-cc-pVTZ. Figure 5(b) shows the C2 basin illustrating its flattened surfaces shared with the C3 and C4 basins. The C3|C4 basins exhibit a DI of 0.986, close to the DI expected for a single bond.

Atom Pairs	$1-C_S$	TS-1→2	$2-C_S$	$3-C_S$
$(\Omega \text{ and } \Omega')$				
C1 C2	1.317 ^a (1.572) ^b	1.460	0.965(1.147)	0.939(1.109)
			(1.147)	
C2 C3	0.590(0.715)	0.934	0.651(0.788)	0.689(0.839)
			(0.788)	
C1 C3	0.087(0.141)	0.085	0.443(0.558)	0.338(0.419)
			(0.557)	
C3 C4	0.986(1.164)	1.051	- ^c	_ c
C1 C4	_ ^c	0.329	_ ^c	_c
C2 C4	_ ^c	0.599	0.066(0.096)	0.065(0.084)
I			(0.096)	
C1 C5	-	-	-	0.886(1.053)

Table 4. Selected delocalization indexes for pairs of atoms at CCSD(full)/6-311+G(2d,p)(tight) and B3PW91/aug-cc-pVTZ(tight).

^a At CCSD(full)/6-311+G(2d,p)(tight). ^b Values in italics in brakets at B3PW91/aug-cc-pVTZ(tight). ^b Values in bold italics in brakets at B3PW91/aug-cc-pVTZ(opt-verytight), int-ultrafine). ^c Data for atom pairs related by symmetry not included.



Figure 5. (a) C1 and C3 basins of $1-C_s$ at CCSD(full)/6-311+G(2d,p)(tight), (b) C2 basin of $1-C_s$ at CCSD(full)/6-311+G(2d,p)(tight).

TS-1→2

Figure 6(a) and (b) display the C1|C4 and C2|C4 basins of **TS-1** \rightarrow 2. Even though there is no BCPs/BPs between these pairs of basins, C1 and C2 are in close proximity to C4. Figure 6(c) shows how C4 is 'distorted' by the presence of C1 and C2. At B3PW91/aug-cc-pVTZ, the C1|C4 and C2|C4 DIs are 0.329 and 0.599 indicating that the delocalization of electrons between C2-C4

is twice as large as it is in the case of C1-C4. Figure 6(d) shows the C3 basin that possesses a 'wedge' of density that intervenes between C2 and C4 and precludes the existence of a BCP/BP.



Figure 6. (a) C1 and C4 basins of **TS-1→2** at B3PW91/aug-cc-pVTZ(tight), (b) C2 and C4 basins of **TS-1→2** at B3PW91/aug-cc-pVTZ(tight), (c) C4 basin of **TS-1→2** at B3PW91/aug-cc-pVTZ(tight), and (d) C3 basin of **TS-1→2** at B3PW91/aug-cc-pVTZ(tight).

Figure 7(a) shows the C1|C3 basins of **2-** C_S at CCSD(full)/6-311+G(2d,p)(tight) that do not exhibit a BCP/BP. *Nevertheless there is a high degree of delocalization of electrons between these basins;* the DI is high at 0.443; 0.588 at B3PW91/aug-cc-pVTZ. Figure 7(b) and (c) display the C1 and C3 basins and clearly show the 'distortion' induced by their proximity. The DI for the C2|C4 pair is 0.066 at CCSD(full)/6-311+G(2d,p)(tight) and 0.096 at B3PW91/aug-cc-pVTZ. We observed this behaviour – relatively large DIs but no BPs – previously for the so-called 7-norbornyl cation¹³ and in the case of trimethylsilyl(carbene) and trimethylgermyl(carbene).¹²



Figure 7. (a), C1 and C3 basins of $2-C_s$ at CCSD(full)/6-311+G(2d,p)(tight), (b) C1 basin of $2-C_s$, (c) C3 basin of $2-C_s$.

Figure 8(a) shows the C1|C3 basins of $3-C_S$ at CCSD(full)/6-311+G(2d,p)(tight) that do not exhibit a BCP/BP. As seen in the case of $2-C_S$ there is a high degree of delocalization of electrons between these basins; the DI is quite large (0.338) but smaller than in the case of $2-C_S$; at B3PW91/aug-cc-pVTZ the value is 0.419. Figure 7(b) and (c) display the C1 and C3 basins that clearly show how they impinge on each other their surfaces are flattened. The DI for the C2|C4 pair is 0.065 at CCSD(full)/6-311+G(2d,p)(tight) and 0.084 at B3PW91/aug-cc-pVTZ.



Figure 8. (a) C1 and C3 basins of $3-C_S$ at CCSD(full)/6-311+G(2d,p)(tight), (b) C1 basin of $3-C_S$, (c) C3 basin of $3-C_S$.

Conclusions

This study shows that the so-called nonclassical bicyclobutonium cation exhibits the molecular structure/graph of a distorted cyclobutyl cation at its equilibrium geometry. It documents another successful application of the QTAIM-DI-VISAB method in establishing the true nature of the bonding in hyper-coordinated so-called nonclassical carbocations; this approach obviates the

need for using arbitrary dotted-line or solid-line representations of bonding in hypercoordinate species.

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