

EXACT TRAVELLING WAVE SOLUTIONS TO THE GENERALIZED KURAMOTO-SIVASHINSKY EQUATION

Changpin Li[†], Guanrong Chen[‡] and Suchuan Zhao[§]

[†]*Department of Mathematics, Shanghai University, Shanghai 200436, P. R. China
leecp@online.sh.cn*

[‡]*Department of Electronic Engineering, City University of Hong Kong, Hong Kong, P. R. China
gchen@ee.cityu.edu.hk*

[§]*Department of Physics, Shanghai University, Shanghai 200436, P. R. China*

Abstract— By using a special transformation, the new exact travelling wave solutions to the generalized Kuramoto- Sivashinsky equation are obtained.

Keywords— travelling wave solutions, solitary wave solutions, Kuramoto-Sivashinsky equation.

I. INTRODUCTION

In this paper, we consider the generalized Kuramoto-Sivashinsky equation (Yang, 1994):

$$u_t + \beta u^\alpha u_x + \gamma u^\tau u_{xx} + \delta u_{xxxx} = 0, \quad (1)$$

where $\alpha, \beta, \gamma, \delta, \tau \in R$ and $\alpha\beta\gamma\delta \neq 0$.

When $\alpha = \beta = 1$ and $\tau = 0$, (1) reduces to the original Kuramoto-Sivashinsky (K-S) equation. The K-S equation was derived by Kuramoto (1978) for the study of phase turbulence in the Belousov-Zhabotinsky reaction. An extension of this equation to two or more spatial dimensions was then given by Sivashinsky (1977, 1980) in the study of the propagation of a flame front for the case of mild combustion. The K-S equation represents one class of pattern formation equation (Yang, 1994; Temam, 1988), and it also serves as a good model of bifurcation and chaos (Abdel-Gawad & Abdusalam, 2001; Li and Chen 2001, 2002).

As far as the travelling wave solutions are concerned, one can always use the transform

$$u(x, t) = u(\xi), \quad \xi = x - ct, \quad (2)$$

where c is the wave velocity. The travelling wave solutions of (1) satisfy the following ordinary differential equation:

$$-cu' + \beta u^\alpha u' + \gamma u^\tau u'' + \delta u'''' = 0. \quad (3)$$

In (Yang, 1994), using the ansatz (Bernoulli equation)

$$u' = au + bu^n, \quad (4)$$

where $a, b, n \in R, ab < 0$ and $n \neq 1$, the exact travelling wave solution to (1) for $\alpha = 3\tau = 9$ was obtained. In this presentation, we further introduce the following ansatz:

$$u(\xi) = v^h(\xi), \quad v' = av + bv^n, \quad (5)$$

where $abh \neq 0, n \neq 1$ and $ab < 0$, and obtain a new exact solution for the equation.

From (5), one first gets

$$u(\xi) = \left[-\frac{a}{2b} \tanh\left(\frac{n-1}{2}a(\xi - c_0)\right) - \frac{a}{2b} \right]^{\frac{h}{n-1}}, \quad (6)$$

in which c_0 is an arbitrary constant. If $h/(n-1) > 0$, (6) is the solitary wave solution connecting the two stationary states $u = 0$ and $u = (-\frac{a}{b})^{h/(n-1)}$ (Lu, et al., 1993). So, the relative orbit is a heteroclinic orbit.

Repeating some differential calculations, one can obtain the following formulas:

$$v'' = (a + nbv^{n-1})v', \quad (7)$$

$$v'''' = [a^3 + a^2bn(n^2 + n + 1)v^{n-1} + 3ab^2n^2(2n-1)v^{2n-2} + b^3n(2n-1)(3n-2)v^{3n-3}]v'. \quad (8)$$

$$u' = hv^{h-1}v', \quad (9)$$

$$u'' = [h^2av^{h-1} + hb(n+h-1)v^{n+h-2}]v', \quad (10)$$

$$u'''' = \{h^4a^3v^{h-1} + ha^2b(n+h-1)[h^2 + (n+h-1)(n+2h-1)]v^{n+h-2} + 3hab^2(n+h-1)^2(2n+h-2)v^{2n+h-3} + hb^3(n+h-1)(2n+h-2)(3n+h-3)v^{3n+h-4}\}v'. \quad (11)$$

Then, by substituting the first formula of (5) and (9)-(11) into (3), one has

$$\begin{aligned} & \{(-ch + \delta h^4 a^3)v^{h-1} + \beta h v^{\alpha h+h-1} + \gamma h^2 a v^{\tau h+h-1} \\ & + \gamma h b(n+h-1)v^{n+(\tau+1)h-2} + \delta h a^2 b(n+h-1) \cdot \\ & [3h^2 + 3h(n-1) + (n-1)^2]v^{n+h-2} + 3\delta h a b^2(n+h-1)^2 \cdot \\ & (2n+h-2)v^{2n+h-3} + \delta h b^3(n+h-1)(2n+h-2) \cdot \end{aligned}$$

$$(3n + h - 3)v^{3n+h-4} \} v' = 0. \tag{12}$$

By furthermore comparing the same orders of v , one can determine values of the parameters a, b, n and h . However, to consider all possible cases is rather complicated. In order to keep the presentation short, only the following interesting cases with $h = 1, n = 0$ and $n = 2$ are considered here.

II. CASE h=1

If $h = 1$, then $u(\xi) = v(\xi)$ and $u' = au + bu^n$, so that (12) is reduced to

$$\begin{aligned} &(\delta a^3 - c) + \beta u^\alpha + \gamma a u^\tau + \gamma b n u^{\tau+n-1} + \\ &\delta a^2 b n (n^2 + n + 1) u^{n-1} + 3 \delta a b^2 n^2 (2n - 1) u^{2n-2} + \\ &\delta b^3 n (2n - 1) (3n - 2) u^{3n-3} = 0. \end{aligned} \tag{13}$$

By comparing the same orders of u , one finds the following situations.

(1) $n = \frac{1}{2}$

(1a) $\tau = 0, \alpha = -\frac{1}{2}$

$$\delta a^3 - c + \gamma a = 0, \beta + \frac{1}{2} \gamma b + \frac{7}{8} \delta a^2 b = 0,$$

that is,

$$b = -\frac{8\beta}{4\gamma + 7\delta a^2}, \quad c = \delta a^3 + \gamma a. \tag{14}$$

In (14), a is a parameter. One should choose a such that $ab < 0$. The same should be done for the similar cases below.

(1b) $\tau = -\frac{1}{2}, \alpha = -1$

$$\delta a^3 - c = 0, \quad \beta + \frac{1}{2} \gamma b = 0, \quad \gamma a + \frac{7}{8} \delta a^2 b = 0,$$

that is,

$$a = \frac{4\gamma^2}{7\delta\beta}, \quad b = -\frac{2\beta}{\gamma}, \quad c = \delta a^3. \tag{15}$$

(2) $n = \frac{2}{3}$

If $n = \frac{2}{3}$, then (13) can be translated into

$$\begin{aligned} &(\delta a^3 - c) + \beta u^\alpha + \gamma a u^\tau + \frac{2}{3} \gamma b u^{\tau-\frac{1}{3}} \\ &+ \frac{38}{27} \delta a^2 b u^{-\frac{1}{3}} + \frac{4}{9} \delta a b^2 u^{-\frac{2}{3}} = 0. \end{aligned}$$

The following results are immediate.

(2a) $\tau = 0, \alpha = -\frac{2}{3}$

$$\delta a^3 - c + \gamma a = 0, \quad \frac{2}{3} \gamma b + \frac{38}{27} \delta a^2 b = 0, \quad \beta + \frac{4}{9} \delta a b^2 = 0.$$

Therefore, one can easily find that

$$\begin{aligned} a &= \pm 3 \left(-\frac{\gamma}{19\delta} \right)^{\frac{1}{2}} (\gamma\delta < 0), \quad b = \mp \frac{3}{2} \left(-\frac{\beta}{\delta a} \right)^{\frac{1}{2}} \\ &(\delta\beta a < 0), \quad c = \delta a^3 + \gamma a. \end{aligned} \tag{16}$$

(2b) $\tau = -\frac{1}{3}, \alpha = -\frac{1}{3}$

$$\delta a^3 - c = 0, \beta + \gamma a + \frac{38}{27} \delta a^2 b = 0, \quad \frac{2}{3} \gamma b + \frac{4}{9} \delta a b^2 = 0,$$

so,

$$a = \frac{9\beta}{10\gamma}, \quad b = -\frac{5\gamma^2}{3\delta\beta}, \quad c = \delta a^3. \tag{17}$$

(2c) $\tau = -\frac{1}{3}, \alpha = -\frac{2}{3}$

$$\delta a^3 - c = 0, \quad \gamma a + \frac{38}{27} \delta a^2 b = 0, \quad \beta + \frac{2}{3} \gamma b + \frac{4}{9} \delta a b^2 = 0,$$

so,

$$a = \frac{90\gamma^2}{361\beta\delta}, \quad b = -\frac{57\beta}{20\gamma}, \quad c = \delta a^3. \tag{18}$$

Besides $n = \frac{1}{2}$ and $n = \frac{2}{3}$, one has the following case.

(3) $\tau = n - 1, \alpha = 3n - 3$

For this case,

$$\begin{aligned} \delta a^3 - c = 0, \quad \gamma a + \delta a^2 b n (n^2 + n + 1) = 0, \\ \gamma b n + 3 \delta a b^2 n^2 (2n - 1) = 0, \quad \beta + \delta b^3 n (2n - 1) (3n - 2) = 0 \end{aligned}$$

From the second and the third equations, it follows that $n = 4$. So, if and only if $n = 4, \alpha = 3\tau = 9$, there exist real number solutions for a, b and c , as

$$a = \frac{\gamma}{6} \left(\frac{5}{49\beta\delta^2} \right)^{\frac{1}{3}}, \quad b = \left(-\frac{\beta}{280\delta} \right)^{\frac{1}{3}}, \quad c = \frac{5\gamma^3}{10584\beta\delta}. \tag{19}$$

This result is the same as that obtained in Yang (1994).

III. CASE n=0

For $n = 0$, (12) is reduced to

$$\begin{aligned} &(\delta a^3 h^3 - c) v^{h-1} + \beta v^{ah+h-1} + \gamma a h v^{\tau h+h-1} + \gamma b (h-1) \cdot \\ &v^{\tau h+h-2} + \delta a^2 b (h-1) (3h^2 - 3h + 1) v^{h-2} + 3 \delta a b^2 (h-1)^2 \cdot \\ &(h-2) v^{h-3} + \delta b^3 (h-1) (h-2) (h-3) v^{h-4} = 0. \end{aligned} \tag{20}$$

After considering the coefficients of some orders of v , one has the following cases, with $h = 1, h = 2$ and $h = 3$, respectively.

(1) $h = 1$

There exists one and only one sub-case with $\alpha = \tau \neq 0$ for $h = 1$ (Note: $\alpha \neq 0$), as follows.

$$(1a) \alpha = \tau \neq 0,$$

$$a = -\frac{\beta}{\gamma}, \quad c = \delta a^3. \quad (21)$$

$$(2) h = 2$$

For $h = 2$, one also has a sub-case.

$$(2a) \tau = 0, \alpha = -\frac{1}{2}$$

For this sub-case, one has

$$7\delta a^2 b + \gamma b + \beta = 0, \quad 8\delta a^3 - c + 2\gamma a = 0.$$

Thus,

$$b = -\frac{\beta}{7\delta a^2 + \gamma}, \quad c = 8\delta a^3 + 2\gamma a. \quad (22)$$

$$(3) h = 3$$

For $h = 3$, (20) can be changed to

$$(27\delta a^3 - c)v^2 + \beta v^{3\alpha+2} + 3\gamma a v^{3\tau+2} + 2\gamma b v^{3\tau+1} + 38\delta a^2 b v + 12\delta a b^2 = 0, \quad (23)$$

and only three cases exist, as follows.

$$(3a) \tau = -\frac{1}{3}, \alpha = -\frac{2}{3}$$

Here,

$$27\delta a^3 - c = 0, \quad 3\gamma a + 38\delta a^2 b = 0, \quad \beta + 2\gamma b + 12\delta a b^2 = 0.$$

Therefore,

$$a = \frac{30\gamma^2}{361\delta\beta}, \quad b = -\frac{19\beta}{20\gamma}, \quad c = 27\delta a^3. \quad (24)$$

$$(3b) \tau = -\frac{1}{3}, \alpha = -\frac{1}{3}$$

It is clear that

$$27\delta a^3 - c = 0, \quad \beta + 3\gamma a + 38\delta a^2 b = 0, \quad 2\gamma b + 12\delta a b^2 = 0.$$

Hence,

$$a = \frac{3\beta}{10\gamma}, \quad b = -\frac{5\gamma^2}{9\delta\beta}, \quad c = 27\delta a^3. \quad (25)$$

$$(3c) \tau = 0, \alpha = -\frac{2}{3}$$

By the same reasoning, one has

$$27\delta a^3 - c + 3\gamma a = 0, \quad 2\gamma b + 38\delta a^2 b = 0, \quad \beta + 12\delta a b^2 = 0,$$

so a, b, c are as below:

$$a = \pm \left(-\frac{\gamma}{19\delta}\right)^{\frac{1}{2}}, \quad \left(\frac{\gamma}{\delta} < 0\right), \quad b = \mp \left(-\frac{\beta}{12\delta a}\right)^{\frac{1}{2}}, \quad \left(\frac{\beta}{\delta a} < 0\right), \quad c = 27\delta a^3 + 3\gamma a. \quad (26)$$

For $h \neq 1$, $h \neq 2$, and $h \neq 3$, a comparison between the corresponding terms in (20) gives only one case, as follows.

$$(4) \alpha = 3\tau = -\frac{3}{h}.$$

It follows that

$$\delta a^3 h^3 - c = 0, \quad \delta a^2 b(h-1)(3h^2 - 3h + 1) + \gamma a h = 0,$$

$$3\delta a b^2(h-1)^2(h-2) + \gamma b(h-1) = 0,$$

$$\delta b^3(h-1)(h-2)(h-3) + \beta = 0.$$

The second and the third equations of the above system give $h = -\frac{1}{3}$. So, if and only if $h = -\frac{1}{3}$, $\alpha = 3\tau = 9$, the above system has real number solutions for a, b and c , as

$$a = -\frac{\gamma}{14\delta} \left(\frac{35\delta}{\beta}\right)^{\frac{1}{3}}, \quad b = \frac{3}{2} \left(\frac{\beta}{35\delta}\right)^{\frac{1}{3}}, \quad c = -\frac{\delta a^3}{27}. \quad (27)$$

IV. CASE n=2

Substituting $n = 2$ into (12) gives

$$(-c + \delta h^3 a^3) v^{h-1} + \delta a^2 b(h+1)(3h^2 + 3h + 1) v^h + 3\delta a b^2(h+1)^2(h+2) v^{h+1} + \delta b^3(h+1)(h+2)(h+3) v^{h+2} + \gamma h a v^{\tau h+h-1} + \gamma b(h+1) v^{\tau h+h} + \beta v^{\alpha h+h-1} = 0. \quad (28)$$

$$(1) h = -1$$

For $h = -1$, there exist only one sub-case.

$$(1a) \tau = \alpha \neq 0$$

$$a = \frac{\beta}{\gamma}, \quad c = -\delta a^3. \quad (29)$$

$$(2) h = -2$$

For $h = -2$, there exist two sub-cases.

$$(2a) \tau = 0, \alpha = -\frac{1}{2}$$

By the same reason, one has

$$-c - 8\delta a^3 - 2\gamma a = 0, \quad -7\delta a^2 b - \gamma b + \beta = 0,$$

so,

$$b = \frac{\beta}{7\delta a^2 + \gamma}, \quad c = -2\gamma a - 8\delta a^3. \quad (30)$$

$$(2b) \tau = -\frac{1}{2}, \alpha = -1$$

Similarly, one has

$$-c - 8\delta a^3 = 0, \quad -7\delta a^2 b - 2\gamma a = 0, \quad -\gamma b + \beta = 0,$$

so,

$$a = -\frac{2\gamma^2}{7\delta\beta}, \quad b = \frac{\beta}{\gamma}, \quad c = -8\delta a^3. \quad (31)$$

$$(3) h = -3$$

For $h = -3$, there exist three sub-cases.

$$(3a) \tau = 0, \alpha = -\frac{2}{3}$$

For this sub-case,

$$c + 27\delta a^3 + 3\gamma a = 0, 38\delta a^2 b + 2\gamma b = 0, 12\delta ab^2 - \beta = 0.$$

The solutions are

$$a = \pm \left(\frac{-\gamma}{19\delta}\right)^{\frac{1}{2}} \left(\frac{\gamma}{\delta} < 0\right), b = \mp \left(\frac{\beta}{12\delta a}\right)^{\frac{1}{2}} \left(\frac{\beta}{\delta a} < 0\right),$$

$$c = -27\delta a^3 - 3\gamma a. \tag{32}$$

$$(3b) \tau = \alpha = -\frac{1}{3}$$

Similarly,

$$c + 27\delta a^3 = 0, 38\delta a^2 b + 3\gamma a - \beta = 0, 12\delta ab^2 + 2\gamma b = 0.$$

The solutions are

$$a = -\frac{3\beta}{10\gamma}, b = \frac{5\gamma^2}{9\delta\beta}, c = -27\delta a^3. \tag{33}$$

$$(3c) \tau = -\frac{1}{3}, \alpha = -\frac{2}{3}$$

The following system is determined by the same reasoning:

$$c + 27\delta a^3 = 0, 38\delta a^2 b + 3\gamma a = 0, 12\delta ab^2 + 2\gamma b - \beta = 0,$$

so,

$$a = -\frac{30\gamma^2}{361\delta\beta}, b = \frac{19\beta}{20\gamma}, c = -27\delta a^3. \tag{34}$$

Besides $h = -1$, $h = -2$ and $h = -3$, there exists only one case, with $h = \frac{1}{\tau}, \alpha = 3\tau$.

$$(4) h = \frac{1}{\tau}, \alpha = 3\tau.$$

It follows that

$$-c + \delta h^3 a^3 = 0, \delta a^2 b(h+1)(3h^2 + 3h + 1) + \gamma ha = 0,$$

$$3\delta ab^2(h+1)^2(h+2) + \gamma b(h+1) = 0,$$

$$\delta b^3(h+1)(h+2)(h+3) + \beta = 0.$$

The second and third equations in the above system give $h = \frac{1}{3}$. So, if and only if $h = \frac{1}{3}$ and $\alpha = 3\tau = 9$, parameters a, b and c are given by

$$a = \frac{\gamma}{2} \left(\frac{5}{49\beta\delta^2}\right)^{\frac{1}{3}}, b = -3 \left(\frac{\beta}{280\delta}\right)^{\frac{1}{3}}, c = \frac{5\gamma^3}{10584\beta\delta}. \tag{35}$$

ACKNOWLEDGEMENTS

This research was supported by the NSF of China (Grant No. 19971057) and the Hong Kong CERG (Grant No. 9040579).

REFERENCES

Abdel-Gawad, H. I. & Abdusalam, H. A., "Approximate solutions of the Kuramoto-Sivashinsky equation for periodic boundary value problems and chaos," *Chaos, Solitons and Fractals*, **12**, 2039-2050 (2001).

Kuramoto, Y., "Diffusion-induced chaos in reaction system," *Prog. Theor. Phys. Suppl.*, **64**, 346-367 (1978).

Li, C. P. & Chen, G., "Bifurcation analysis of the Kuramoto-Sivashinsky equation in one spatial dimension," *Int. J. Bifurcation and Chaos*, **11** (9), 2493-2499 (2001).

Li, C. P. & Chen, G., "Bifurcation from an equilibrium of the steady state Kuramoto-Sivashinsky equation in two spatial dimensions," *Int. J. Bifurcation and Chaos*, **12** (1), 103-114 (2002).

Lu, B. Q., Xiu, B. Z., Pang, Z. L. & Jiang, X. F., "Exact traveling wave solutions of one class of nonlinear diffusion equations," *Phys. Lett. A*, **175**, 113-115 (1993).

Sivashinsky, G. I., "Nonlinear analysis of hydrodynamic instability in laminar flames, Part I, derivation of basic equations," *Acta Astronautica*, **4**, 1177-1206 (1977).

Sivashinsky, G. I., "On flame propagation under conditions of stoichiometry," *SIAM J. Appl. Math.*, **39**, 67-82 (1980).

Temam, R., *Infinite Dimensional Dynamical Systems in Mechanics and Physics*, Springer-Verlag, New York (1988).

Yang, Z. J., "Travelling wave solutions to nonlinear evolution and wave equations," *J. Phys. A: Math. Gen.*, **27**, 2837-2855 (1994).

Received: April 30, 2002.
Accepted for publication: July 25, 2002.
Recommended by Subject Editor Jorge L. Moiola.