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**Impact of diurnal  
variability in SST on  
CO<sub>2</sub> flux**

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# The impact of diurnal variability in sea surface temperature on the atlantic air-sea CO<sub>2</sub> flux

H. Kettle<sup>1</sup>, C. J. Merchant<sup>1</sup>, C. D. Jeffery<sup>2</sup>, M. J. Filipiak<sup>1</sup>, and C. L. Gentemann<sup>3</sup>

<sup>1</sup>School of GeoSciences, The University of Edinburgh, UK

<sup>2</sup>National Oceanography Centre, Southampton, UK

<sup>3</sup>Remote Sensing Systems, Santa Rosa, CA, USA

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Correspondence to: H. Kettle (h.kettle@ed.ac.uk)

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## Abstract

The effect of diurnal variations in sea surface temperature (SST) on the air-sea flux of CO<sub>2</sub> over the central Atlantic ocean and Mediterranean Sea is evaluated for 2005–2006. We use high resolution hourly satellite SST data to determine the diurnal warming (ΔSST). The CO<sub>2</sub> flux is then computed using three different temperature fields – a foundation temperature ( $T_f$ , measured at a depth where there is no diurnal variation),  $T_f$  plus the hourly ΔSST and  $T_f$  plus monthly-averaged ΔSST. This is done in conjunction with a physically-based parameterisation for the gas transfer velocity (NOAA-COARE). The differences between the fluxes evaluated for these three different temperature fields quantifies the effects of both diurnal warming and diurnal covariations. We find that including diurnal warming increases the CO<sub>2</sub> flux out of the Atlantic for 2005–2006 from 9.6 Tg C a<sup>-1</sup> to 30.4 Tg C a<sup>-1</sup> (hourly ΔSST) and 31.2 Tg C a<sup>-1</sup> (monthly ΔSST). Diurnal warming, therefore, has a large impact on the annual net CO<sub>2</sub> flux but diurnal covariations in variables are negligible implying that CO<sub>2</sub> fluxes may be adequately computed using monthly averaged ΔSSTs along with a suitable foundation temperature.

## 1 Introduction

During the day, the upper 2 m of the ocean typically absorbs about 50% of the solar radiation reaching its surface. At night this layer then cools, losing heat to the atmosphere through radiative latent and sensible heat fluxes. This diurnal heating and cooling can lead to significant variations in the sea surface temperature (SST) (e.g., Stuart-Menteth et al., 2003; Gentemann et al., 2003). Here we investigate the impact of diurnal variability in SST on CO<sub>2</sub> fluxes by using SST data from the Spinning Enhanced Visible and Infrared Imager (SEVIRI) geostationary satellite. Typical regional and seasonal variations in diurnal warming over the SEVIRI disk region are shown in Fig. 1. Due to averaging, diurnal changes in SST (ΔSST) shown in Fig. 1 are only up to 1.5 K but on

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individual days localised warming can be as much as 6 K within a shallow 'warm layer' at the sea surface (Merchant et al., 2008; Stramma et al., 1986).

The sea-air flux of CO<sub>2</sub>,  $F$ , is controlled by the transfer of CO<sub>2</sub> across the aqueous boundary layer, such that,

$$F = k([\text{CO}_{2w}] - [\text{CO}_{2s}]) \quad (1)$$

(McGillis and Wanninkhof, 2006) where  $k$  is the gas transfer velocity,  $[\text{CO}_{2w}]$  is the CO<sub>2</sub> concentration at the base of the mass boundary layer and  $[\text{CO}_{2s}]$  is the CO<sub>2</sub> concentration at the surface skin. Because dissolved CO<sub>2</sub> in the ocean is strongly buffered by dissolved inorganic carbon species, the transfer of CO<sub>2</sub> across the interface does not significantly affect the total dissolved CO<sub>2</sub> concentration (i.e., we assume  $[\text{CO}_{2w}]$  is not affected by the flux). The concentration of CO<sub>2</sub> may be expressed as a combination of the solubility of CO<sub>2</sub> in sea water and its partial pressure, so that Eq. 1 becomes:

$$F = k(\alpha_w p\text{CO}_{2w} - \alpha_s p\text{CO}_{2a}) \quad (2)$$

where  $\alpha$  is the solubility of CO<sub>2</sub> in sea water,  $p\text{CO}_2$  is the partial pressure of CO<sub>2</sub> and the subscripts  $w$ ,  $s$  and  $a$  denote the bottom of the mass boundary layer, the skin and the air respectively. Note that  $p\text{CO}_2$  is assumed to be equivalent to the fugacity of CO<sub>2</sub> (<0.5% error over the relevant temperature range; McGillis and Wanninkhof, 2006). Since each of these factors ( $k$ ,  $p\text{CO}_2$  and  $\alpha$ ) vary with SST (Fig. 2),  $F$  will also vary diurnally. Thus, this study uses high resolution satellite measurements of the ocean skin to estimate CO<sub>2</sub> flux.

Previous studies have suggested the "thermal skin effect" (cooling/warming of the upper few millimetres of the ocean) affects flux (e.g. Robertson and Watson, 1992; Van Scoy et al., 1995), as does the warming of the upper few metres of the ocean by solar radiation (McNeil and Merlivat, 1996). Work by Olsen et al. (2004) and McNeil and Merlivat (1996) on this topic differs from the study herein, in that they use a wind-based parameterisation for the gas transfer velocity and averaged values of

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diurnal warming. These two simplifications may underestimate the importance of diurnal warming. This is because averaging eliminates covariations between variables and wind-based transfer velocities predict no gas flux when there is no wind, which are the conditions under which large  $\Delta$  SSTs may occur. Moreover, a field experiment has shown that it is possible for CO<sub>2</sub> fluxes to have only a weak dependence on wind speed but a strong dependence on the diurnal heating cycle (e.g., GasEx-2001 in the Equatorial Pacific; McGillis et al., 2004). Therefore, in this study we use a more complex physically-based parameterisation that includes buoyancy driven, as well as wind driven, gas transfer, by Fairall et al. (2000) with modifications by Jeffery et al. (2007); along with a slightly different formulation for the flux (Eq. 2) as recommended by McGillis and Wanninkhof (2006).

Since we are investigating gas flux through the air-sea interface, we define the SST to be the temperature of the ocean skin ( $T_s$ ). We can express this in terms of a foundation (or bulk) temperature (Donlon et al., 2007) below the diurnally warmed layer ( $T_f$ ), the temperature difference associated with diurnal heating ( $\Delta T_{dw}$ ) and the temperature drop across the skin, ( $\Delta T_s$ ) such that,

$$T_s = T_f + \Delta T_{dw} - \Delta T_s. \quad (3)$$

The impact of temperature on CO<sub>2</sub> flux is investigated by computing fluxes over the SEVIRI disk region for the following three scenarios:

**Scenario 1.** The specified temperature is equivalent to the foundation temperature. This is the most commonly used temperature and ignores both the skin effect and diurnal variations.

$$T(x, y, t) = T_f(x, y, t), \quad (4)$$

We use this to calculate CO<sub>2</sub> flux over the complete SEVIRI disk with no account taken of diurnal warming – the flux generated (using Eq. 2) is denoted  $F_f$ .

**Scenario 2.** This examines the effect of diurnal variability on CO<sub>2</sub> flux by using a temperature with hourly estimates of the diurnal variability included (Eq. 3).

$$T(x, y, t) = T_s(x, y, t) \quad (5)$$

$$= T_f + \Delta T_{dw} - \Delta T_s.$$

For a given time slot this will only cover a small fraction of the SEVIRI disk because we only use locations where warming occurs and where there are data available (i.e. pixels not obscured by cloud).  $\Delta T_{dw} - \Delta T_s$  is estimated from satellite measurements (see Sect. 3.1) and  $T_f$  is as in Scenario 1. The flux computed under these conditions is denoted  $F_{dv}$ .

**Scenario 3.** This investigates the impact of using a monthly average diurnal variability, instead of hourly estimates, on CO<sub>2</sub> flux. The temperature used is the foundation temperature plus the monthly mean warming, (denoted by  $\overline{\Delta T_{dw} - \Delta T_s}$ ):

$$T(x, y, t) = T_f(x, y, t) + \overline{\Delta T_{dw}(x, y) - \Delta T_s} \quad (6)$$

This gives flux estimates accounting for the warming but not the time variability, denoted  $F_w$

The difference between the CO<sub>2</sub> flux fields resulting from scenarios 1 and 2 ( $F_{dv} - F_f$ ) examines the effect of the increase in SST caused by diurnal warming. The difference between scenarios 2 and 3 ( $F_{dv} - F_w$ ) examines the effect of the covariability of SST with the other factors affecting flux.

## 2 Data

Satellite observations of SST, surface solar irradiance (SSI) and downward longwave irradiance (DLI) are provided by EUMETSAT's Ocean and Sea-Ice Satellite Application Facility (OSISAF), and consist of hourly fields over a field of view that encompasses the east Atlantic Ocean and the Mediterranean Sea (Fig. 1). SSTs are derived from the SEVIRI radiances (OSISAF, Atlantic Sea Surface Temperature Product Manual, Version 1.6, October 2006, [http://www.osi-saf.org/biblio/docs/ss1\\_pmatlst\\_1.6.pdf](http://www.osi-saf.org/biblio/docs/ss1_pmatlst_1.6.pdf)). The resolution of the data is 0.05° and geographical coverage is 60° S to 60° N, 60° W to 45° E (the disk is approximately a fifth of Earth's total surface area). Data with satellite zenith angle greater than 60° were excluded due to the potential unreliability of cloud

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screening and poorer SST precision. The difference between the SEVIRI SSTs and matched drifting buoys (between July 2004 and July 2005) has mean standard deviation of  $0.01 \pm 0.49$  K which includes both drifter errors and spatially correlated retrieval errors. SSTs are only measurable when the sky is clear, so each data point is assigned a confidence level ranging from 1 (“bad”) to 5 (“excellent”), depending on the possible cloud contamination (LeBorgne et al., 2006). We bin the data onto a  $0.2^\circ$  grid to increase the apparent completeness (in space and time) of the SST data and decreases the SST error in a cell due to retrieval noise. This spatial averaging may dampen the amplitudes of very localised diurnal warming but was necessary due to computing constraints. This SST dataset is used to compute the diurnal warming,  $\Delta T_{dw} - \Delta T_s$  (see Methods section).

In addition to the SST dataset described above, a foundation SST data set, provided by Meteo-France, is also used. This is an analysis of night-time sub-skin SSTs optimally interpolated to 00:00 UTC daily. It is this dataset that is used to provide the values of  $T_f$  in the three scenarios (Eq. 4–6).

The wind speeds used in this analysis are the NASA Atlas First-Look (FLK) version 1.1 derived surface winds level 3.0 product which uses available passive microwave satellite wind speeds produced by Remote Sensing Systems and described at <http://sivo.gsfc.nasa.gov/oceanwinds/>. All satellite measurements are processed in a consistent manner using a physically-based retrieval algorithm to determine the wind speed (Wentz, 1997). These wind speeds are used to derive a global 10-m wind speed every 6 h on a 25 km grid using variational analysis method (VAM). These data were linearly interpolated in time and space onto the hourly SEVIRI  $0.05^\circ$  grid. Finally the wind speed data coincident with the grid points of the  $0.2^\circ$  grid used in this study are extracted.

Other meteorological data, pressure ( $P$ ), dew point temperature ( $T_{dew}$ ) and air temperature ( $T_{air}$ ), are taken from the ECMWF operational dataset (N80 Gaussian gridded analysis on surface levels; in ERA-40 format) at 6-hourly intervals) and we linearly interpolate these in time and space.  $\rho CO_{2w}$  and salinity ( $S$ ) are taken from Taka-

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hashi's climatology (Takahashi et al., 2002) – hereafter referred to as Taka02 – which is monthly and on 4° lon×5° lat grid. Where the Taka02 data are not fully resolved we interpolate longitudinally. We use a monthly climatological dataset for the mixed layer depth (MLD) obtained from Scripps Institution of Oceanography (available from <http://ingrid.ldgo.columbia.edu/SOURCES/.IGOSS/.sio>) with 5°lon×2°lat resolution.

### 3 Methods

#### 3.1 Deriving the diurnal variations in SST

The SEVIRI satellite measures  $T_s$  but the processed dataset is corrected for the cool skin by adding 0.2 K. We reverse this correction to retain the original  $T_s$  measurement. To calculate the diurnal warming, at each hour where there is a SST measurement with confidence level 5, we compute the difference between it and the “satellite foundation temperature” ( $T_{sf}$ ) which we define to be the satellite measured temperature just before the time of local dawn ( $t_d$ ). Note this is not the same as the foundation temperature previously mentioned ( $T_f$ ) which is from a different dataset.  $T_{sf}$  throughout the rest of the day is approximated using a linear interpolation between consecutive pre-dawn temperatures, such that

$$T_{sf}(t) = T_s(t_d) + \frac{(T_s(t_d + 24) - T_s(t_d))}{24}(t - t_d) \quad (7)$$

The diurnal temperature difference at time  $t$  is then given by:

$$\Delta T_{dw}(t) - \Delta T_s(t) = T_s(t) - T_{sf}(t) \quad (8)$$

#### 3.2 Computing the CO<sub>2</sub> flux

The sea-air flux of CO<sub>2</sub> (Eq. 2) contains 3 factors which depend on temperature in different ways (see Fig. 2). The following subsections describe the details of how each of these factors is computed and its reliance on SST.

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### 3.2.1 Solubility, $\alpha$

The solubility,  $\alpha$  of CO<sub>2</sub> in sea water is a physical property that determines how much CO<sub>2</sub> will dissolve. CO<sub>2</sub> is poorly soluble in water and its solubility is highly temperature dependent. Solubility (in mol m<sup>-3</sup> atm) can be calculated according to Weiss (1974) by

$$\alpha = 1000 \exp\left(a_1 + a_2 \frac{100}{T_k} + a_3 \log\left(\frac{T_k}{100}\right) + S\left\{b_1 + b_2 \frac{T_k}{100} + b_3 \left(\frac{T_k}{100}\right)^2\right\}\right) \quad (9)$$

where  $T_k$  is the water temperature (Kelvin),  $a_1 = -58.0931$ ,  $a_2 = 90.5069$ ,  $a_3 = 22.2940$ ,  $b_1 = 0.027766$ ,  $b_2 = -0.02588$ ,  $b_3 = 0.0050578$ . As the temperature increases the solubility decreases, e.g., dropping to 40% of its value for a temperature increase from 5° C to 40° C (Fig. 2). To compute the CO<sub>2</sub> flux (Eq. 2) for the different scenarios we evaluate  $\alpha$  using the foundation temperature ( $\alpha_w$ ) and using the skin temperature ( $\alpha_s$ ).

### 3.2.2 Gas transfer velocity, $k$

The gas transfer velocity describes the rate at which a gas moves between the sea and air. The magnitude of the transfer rate is controlled by the thickness of the boundary layer which is a function of near surface turbulence and diffusion. Thus, the transfer rate is determined by the state of the sea surface: by factors such as wave age, fetch, wind speed, the prevalence of bubbles, boundary layer stability and naturally occurring surfactants (e.g. Woolf, 1997; Monahan and Spillane, 1984; Liss and Merlivat, 1986; Asher and Wanninkhof, 1998). It is highly unlikely, therefore, that only one physical variable can completely determine the spatial scales and environmental conditions necessary to predict  $k$ . Despite this, many empirical relationships for  $k$  in practical use are solely functions of wind speed as this is an influential and easily obtainable parameter. Three commonly used wind-based parameterisations are the piecewise linear relation (Liss

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and Merlivat, 1986), the quadratic relation (Wanninkhof, 1992; Nightingale et al., 2000), and the cubic relation (Wanninkhof and McGillis, 1999).

Using this type of parameterisation to examine the influence of diurnal warming on gas flux will likely result in an under-estimation of the effect because at low wind speeds (when diurnal warming is at its most significant) these parameterisations predict virtually no gas flux. To overcome this limitation we use the NOAA Coupled Ocean Atmosphere Response Experiment (COARE) gas transfer parameterisation (Fairall et al., 2000) which is physically (rather than empirically) based. We also include a modification to this parameterisation by Jeffery et al. (2007) to include the effects of nighttime convective overturn of the water column. A brief description of this method is given below.

Fairall et al. (2000) express the transfer velocity as:

$$k = \left( \frac{r_w}{u_{*w}} + \frac{r_a \alpha_n}{u_{*a}} \right)^{-1}, \quad (10)$$

where  $\alpha_n$  is non-dimensionalised solubility ( $\frac{\alpha R_{\text{gas}}}{T}$ ; where  $R_{\text{gas}}$  is the universal gas constant),  $r$  is the 'resistance' and  $u_*$  is the friction velocity (subscripts  $a$  and  $w$  refer to the air and water sides respectively). The resistances are given by:

$$r_w = h_w S_{cw}^{\frac{1}{2}} + \ln \left( \frac{z_{wr}}{\delta} \right) / \kappa, \quad (11)$$

$$r_a = h_a S_{ca}^{\frac{1}{2}} + C_{da}^{\frac{1}{2}} - 5 + \frac{\ln(S_{ca})}{2K}, \quad (12)$$

where  $S_c$  is Schmidt number,  $z_{wr}$  is the measurement depth,  $\delta$  is the thickness of the cool skin,  $C_{da}$  is the airside drag coefficient and  $\kappa$  is the von Karman constant (0.41). The  $h$  factors are concerned with the transport through the cool skin layer and are given by  $h_a = 13.3$  and  $h_w = \frac{13.3\lambda}{6A}$  (Saunders, 1967; Soloviev and Schlüssel, 1994) where  $\lambda$  is computed according to Fairall et al. (1996a) and  $A$  is a tunable constant ( $\approx 1$ ). If there is no cool skin present  $\lambda$  is set to 6.

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Fairall et al. (2000) define the water-side friction velocity  $u_{*w}$  by

$$u_{*w} = u_{*a} \left( \frac{\rho_w}{\rho_a} \right)^{-\frac{1}{2}}. \quad (13)$$

However, in order to include the increased gas transfer caused by convective overturn, Jeffery et al. (2007) modified the expression for  $u_{*w}$  to include waterside “gustiness”.

5 Thus  $u_{*w}$  is newly defined as

$$u_{*w} = \sqrt{C_{dw} S_w^2} \quad (14)$$

where  $C_{dw}$  is the waterside drag coefficient and  $S_w$  is an average value of “wind speed”, which following Stull (1994) and Godfrey and Beljaars (1991) for the airside, is expressed as

$$10 \quad S_w^2 = u_{\text{ref}}^2 + w_g^2, \quad (15)$$

where  $u_{\text{ref}}$  is analogous to a wind speed at some reference depth ( $z_{\text{ref}}$ ), which we can define as

$$u_{\text{ref}} = \frac{u_{*w}}{K} \ln \left( \frac{z_{\text{ref}}}{z_0} \right). \quad (16)$$

The convective buoyancy/velocity scale,  $w_g$  is defined as

$$15 \quad w_g = \beta (-B_f Z_m)^{\frac{1}{3}}, \quad (17)$$

where  $\beta$  is the (tunable) “gustiness parameter”,  $Z_m$  is the depth of the convective layer (we use monthly climatological MLD) and  $B_f$  is the buoyancy flux given as the sum of the buoyancy caused by heating and that caused by freshening through evaporation, such that

$$20 \quad B_f = \frac{g}{\rho_w} \left( \frac{a_1 Q_{\text{net}}}{C_p} - \frac{b_e Q_{\text{lat}}}{L_v} \right), \quad (18)$$

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where  $g$  is the acceleration due to gravity,  $a_1$  is the thermal expansion coefficient ( $2.1 \times 10^{-5} (T + 3.2)^{0.79} \text{ K}^{-1}$ ),  $b_e$  is the saline expansion coefficient (0.026),  $C_p$  is the thermal heat capacity of water and  $\rho_w$  is the density of seawater (both functions of temperature),  $L_v$  is the latent heat of vaporization ( $(2.501 - 0.00237T) \times 10^6 \text{ J kg}^{-1}$ ),  $Q_{\text{net}}$  is the net heat flux (positive into the water) and  $Q_{\text{lat}}$  is the latent heat of evaporation (positive out of water). When the buoyancy flux is positive  $w_g$  is set to zero as the fluxes serve to stabilize the exchange by adding buoyancy to the surface.

Bubble mediated gas transfer ( $k_b$ ) is accounted for by modifying the gas transfer Eq. (10) as follows:

$$k = k_b + \left( \frac{r_w}{u_{*w}} + \frac{r_a \alpha_n}{u_{*a}} \right)^{-1} \quad (19)$$

where  $k_b$  is defined by Woolf's (1997) parameterisation:

$$k_b = BV_0 f \alpha_n \left[ 1 + (e \alpha_n S_{c_w}^{\frac{1}{2}})^{-\frac{1}{n}} \right]^{-n} \quad (20)$$

where

$$f = 3.84 \times 10^{-6} u_{10}^{3.41} \quad (21)$$

and  $V_0 = 6.8 \times 10^{-3} \text{ m s}^{-1}$ ,  $e = 14$ ,  $n = 1.2$  and  $B$  is a tunable constant.

To solve Eqs. 10 to 21 we first compute the heat fluxes ( $Q_{\text{lat}}$  and  $Q_{\text{net}}$ ), the cool skin parameters ( $\delta$  and  $\lambda$ ) and the drag coefficients ( $C_d$ ) using code from the air-sea matlab toolbox from Woods Hole Science Center (<http://woodhole.er.usgs.gov/operations/sea-mat/index.html>; Fairall et al., 1996a, 2000). This requires relative humidity (function of  $T_{\text{air}}$  and  $T_{\text{dew}}$ ), pressure, air temperature (all from ECMWF), wind speed (satellite data), net short wave radiation (SEVIRI SSI) and net long wave radiation (SEVIRI DLI minus the long wave radiation emitted from the ocean).

The (dimensionless) Schmidt number (used in Eqs. 11 and 12) is the kinematic viscosity of the fluid divided by the molecular diffusion coefficient of the gas. For CO<sub>2</sub>

in seawater  $S_{CW}$  can be estimated from a relationship with temperature (Wanninkhof, 1992) such that

$$S_{CW} = 2073.1 - 125.62T + 3.6276T^2 - 0.043219T^3 \quad (22)$$

where  $T$  is in °C. The Schmidt number for CO<sub>2</sub> in air,  $S_{ca}$  is kept constant at 0.8 (Fairall et al., 2000) and is much smaller than its waterside equivalent (~600) so that the transfer resistance for CO<sub>2</sub> is much greater in water than in air.

The gas transfer parameterisation thus contains three empirical parameters which allow tuning to specific data sets:  $A$  (related to the thermal sublayer),  $B$  (related to bubble mediated transfer) and  $\beta$  (the “gustiness” parameter which is related to convective buoyancy effects). Published values of  $A$  and  $B$  derived from CO<sub>2</sub> air-sea flux field experiments are:  $A=0.625$ ,  $B=2.0$  (GasEx 98 – a warm core eddy; Hare et al., 2004), and  $A=1.3$ ,  $B=0.82$  (GasEx 2001 – in the eastern Pacific south of the upwelling region; derived from results by McGillis et al., 2004). Soloviev and Schüssel (1994) use  $A=1.85$  and  $B=1$  based on radon experiments. Thus, there is a significant amount of uncertainty in these two parameters. The gustiness parameter,  $\beta$ , has published values of 1.25 (Fairall et al., 1996b), 1.0 (Miller et al., 1991) and 0.7 (Schumann, 1988) – but note that these are for air. Here we are not tuning the parameterisation to a particular data set so we take the generic values of  $\beta=1$ ,  $A=1$ , and  $B=1$ , since these are roughly the mean of the previously published values and are a neutral choice with no scaling up or down. The effect of SST on the gas transfer velocity is shown in Fig. 2 for a steady wind speed of 2 m s<sup>-1</sup> and in Fig. 3 for a range of wind speeds.

The temperature used to evaluate  $k$  (Eq. 10–22) is changed according to the three scenarios outlined in the Introduction.

### 3.2.3 The partial pressures of CO<sub>2</sub> in the air and sea, $p\text{CO}_2$

We calculate  $p\text{CO}_{2a}$  and  $p\text{CO}_{2w}$  in Eq. 2 based on changes to the Takak02 climatology caused by short term changes in air pressure and SST. Variations in  $p\text{CO}_{2a}$  with

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change in dry air pressure have been shown to be important for flux calculations by Kettle and Merchant (2005). Thus we compute  $p\text{CO}_{2a}$  from:

$$p\text{CO}_{2a}(t) = [\text{CO}_2]_{\text{air}}(P(t) - \text{SVP}(t)) \quad (23)$$

where  $[\text{CO}_2]_{\text{air}}$  are the zonal mean  $\text{CO}_2$  concentrations in the dry atmosphere for 1995 (Globalview- $\text{CO}_2$ , 2000; and used in Taka02) and SVP is the saturation vapour pressure, which is a function of SST and salinity:

$$\begin{aligned} \text{SVP}(t) = 1013.25 \exp\left(24.4543 - 67.4509 \frac{100}{T_k(t)} \right. \\ \left. - 4.8489 \ln\left(\frac{T_k(t)}{100}\right) - 0.000544 S\right), \end{aligned} \quad (24)$$

where salinity,  $S$ , is taken from Taka02.

The change in  $p\text{CO}_{2w}$  with temperature is given by Takahashi et al. (1993) as

$$\frac{\partial \ln(p\text{CO}_{2w})}{\partial T} = 0.0423. \quad (25)$$

However, it should be noted that  $0.0423 \text{ } ^\circ\text{C}^{-1}$  is an approximation and can range between  $0.037$  to  $0.053 \text{ } ^\circ\text{C}^{-1}$  depending upon the carbonate dissociation constants used (McGillis and Wanninkhof, 2006). We compute  $p\text{CO}_{2w}(t)$  based on changes from a reference field, such that:

$$p\text{CO}_{2w}(t) = p\text{CO}_{2w}^{\text{Tak}} \exp(0.0423(T_f(t) - T^{\text{Tak}})), \quad (26)$$

where  $t$  is time,  $p\text{CO}_{2s}^{\text{Tak}}$  and  $T^{\text{Tak}}$  are  $p\text{CO}_{2s}$  and temperature from Taka02 (monthly). Note that  $p\text{CO}_{2w}$  is always computed using the foundation temperature but  $p\text{CO}_{2a}$  will vary between the three scenarios.

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## 4 Results

Figure 4 shows seasonal averages of the CO<sub>2</sub> flux field for 2005–2006 computed using the analysed foundation SST ( $T_f$ ). There is strong outgassing of around 2 mol CO<sub>2</sub> m<sup>-2</sup> a<sup>-1</sup> from the ocean around the equator, changing to ocean uptake of CO<sub>2</sub> beyond the Tropics towards the North and South poles. Regions around 30–40° N and 30–40° S change seasonally between being sources and sinks of CO<sub>2</sub>. The plots compare well in terms of magnitude and spatial distribution with the mean annual flux for 1995 shown by Taka02 and serve as a check that the more complex physically-based flux parameterisation, with far more variables, is generally equivalent (at these scales) to other methods.

Using the foundation SST ( $F_f$ , scenario 1) gives a mean net mass flux of 9.6 Tg C a<sup>-1</sup> out of the ocean over the SEVIRI disk region for 2005–2006. When satellite-measured diurnal variations are included ( $F_{dv}$ , scenario 2) this is increased to 30.4 Tg C a<sup>-1</sup>, and when diurnal warming is represented by a monthly-averaged value ( $F_w$ , scenario 3) the mass flux is 31.2 Tg C a<sup>-1</sup>. Figure 5 shows how using the three different SST datasets affects the total mass flux over the SEVIRI disk for each month during 2005 and 2006. Using satellite-measured diurnal variations increases outgassing by 21.7 Tg C (2005) and 20.0 Tg C (2006) (Fig. 5b). When time covariations are eliminated by using the monthly averaged  $\Delta$ SST the outgassing is increased (or ingassing is reduced) by a further 0.92 Tg C (2005) and 0.69 Tg C (2006). Diurnal covariations reduce the outgassing flux because  $\Delta$ SST and wind speed (the dominant factor affecting flux) are negatively correlated leading to less flux when  $\Delta$ SST (and  $\alpha_s$ ) is high (due to the low wind speed). However, the difference between  $F_{dv}$  and  $F_w$  is small compared with the difference with  $F_f$  implying that the covariation effects of diurnal variability are much less important than the mean effect of diurnal warming. In Fig. 5c the number of  $\Delta$ SST data points derived from satellite measured SSTs each month is plotted. The number of valid measurements range from a minimum in January 2006 (0.67 million) to a maximum in June 2005 (2.01 million).

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Since the effect of diurnal covariability is small, in the rest of this section we focus on the differences between using  $F_{dv}$  (because it is based on the most detailed data available) and  $F_f$  (because it is the foundation SST which is most commonly used for estimating CO<sub>2</sub> flux). Figure 6 shows the spatial distribution of the mean monthly differences in flux caused by diurnal variability, (i.e.,  $F_{dv} - F_f$ ), averaged over 2005–2006 for each month. Here we see that including diurnal variability in SST either causes an increase in the outgassing of CO<sub>2</sub> from the ocean or no change in flux, everywhere over this region. The maximum increase in monthly-averaged outgassing is  $\sim 0.2$  mol CO<sub>2</sub> m<sup>-2</sup> a<sup>-1</sup>. The seasonal maximum is around 5 mol CO<sub>2</sub> m<sup>-2</sup> a<sup>-1</sup> (Fig. 4), however, large regions have zero net flux so the impact of diurnal variability in SST on flux, is regionally very significant. As expected, the impact changes spatially with time of year, with large increases in the Mediterranean in the northern summer ( $\sim 0.1$  mol CO<sub>2</sub> m<sup>-2</sup> a<sup>-1</sup>) and around South America from June to January. The spatial distribution of the available data points is shown in Fig. 7. The figure shows there are some areas where data is very sparse – this is discussed further in the following section.

## 5 Discussion

The results show that including the increase in SST due to diurnal warming acts to increase the outgassing/reduce the ingassing flux of CO<sub>2</sub> from the ocean over the SEVIRI disk region (all other factors being equal). The main factor in the flux Eq. (2) through which  $\Delta$ SST affects flux can be ascertained by differentiating with respect to temperature:

$$\frac{\partial F}{\partial T} = \left( k\alpha_w \frac{\partial p\text{CO}_{2w}}{\partial T} + \alpha_w p\text{CO}_{2w} \frac{\partial k}{\partial T} + k p\text{CO}_{2w} \frac{\partial \alpha_w}{\partial T} \right) - \left( k\alpha_s \frac{\partial p\text{CO}_{2a}}{\partial T} + \alpha_s p\text{CO}_{2a} \frac{\partial k}{\partial T} + p\text{CO}_{2a} k \frac{\partial \alpha_s}{\partial T} \right). \quad (27)$$

Numbering the terms on the right hand side of Eq. 27 from 1–6; we can ignore terms 1 and 3 since these are evaluated at the foundation temperature and will not be affected by diurnal warming. Term 4 can be assumed negligible since  $p\text{CO}_{2a}$  does not vary much with temperature (Fig. 2) so that Eq. 27 becomes:

$$5 \quad \frac{\partial F}{\partial T} = (\alpha_w p\text{CO}_{2w} - \alpha_s p\text{CO}_{2a}) \frac{\partial k}{\partial T} - p\text{CO}_{2a} k \frac{\partial \alpha_s}{\partial T}. \quad (28)$$

Since the partial pressure of  $\text{CO}_2$  in the ocean and atmosphere is approximately in balance the first term on the right hand side of Eq. 28 is close to zero, implying that the diurnal change in flux is dominated by the change in solubility caused by variations in the ocean skin temperature. Solubility decreases with temperature so this term is  
 10 negative, indicating that the flux in the outgassing direction will be increased by diurnal warming. In other words, the change in flux due to diurnal warming can be estimated very approximately by:

$$\Delta F_{\text{diurnal warming}} \approx -p\text{CO}_{2a} k \Delta \alpha_s. \quad (29)$$

However,  $\text{CO}_2$  flux is not just affected by temperature but also by biological activity. Photosynthesis by phytoplankton removes dissolved inorganic carbon (DIC) from the  
 15 surface waters, lowering  $p\text{CO}_{2w}$  when there is sufficient light and nutrient available. Since light availability also varies diurnally the biological effect, which acts to increase  $\text{CO}_2$  ingassing, may eliminate the diurnal increases in outgassing caused by diurnal warming. We estimate the approximate magnitude of the biological effect as follows: Morel and Antoine (2002) show the average net primary production (NPP) over June  
 20 2001 or December 2000 to have a global maximum of  $2 \text{ gC m}^{-2} \text{ d}^{-1}$  (which incidentally is higher than estimates given by Behrenfeld et al. (2005), and Behrenfeld and Falkowski, 1997). We convert this to the NPP over daylight hours by doubling it (assuming no photosynthesis over night, and assuming this is half of the day), and we  
 25 assume photosynthesis occurs over a depth of 100 m, so that the uptake of carbon from the water in one day is  $40 \text{ mg C m}^{-3}$ . This is equivalent to a decrease in DIC of  $3.3 \mu\text{mol C l}^{-1}$ . Using equations representing the sea-water acid-base system with

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expressions for the dissociation constants of carbonic acid, hydrogen carbonate, boric acid and water from DOE (1994) we can compute the change in  $p\text{CO}_{2w}$  for a given change in DIC. If DIC decreases by  $3.3 \mu\text{mol C l}^{-1}$ , and assuming a standard DIC concentration of  $2058 \mu\text{mol l}^{-1}$ , alkalinity of  $2396 \mu\text{mol l}^{-1}$  (e.g. Palmer and Totterdell, 2001) and  $\text{SST}=25^\circ \text{C}$ , the decrease in  $p\text{CO}_{2w}$  is  $5.7 \mu\text{atm}$ . This would result in a change in flux ( $\Delta F_{\text{bio}}$ ) of  $5.7 \times 10^{-6} k \alpha_w$ , which at  $25^\circ\text{C}$  is  $\sim 1.65 \times 10^{-4} k$ . Equating this with Eq. 29 (and assuming  $p\text{CO}_{2a} \approx 350 \mu\text{atm}$ ), indicates that a change in solubility of  $0.47 \text{ mol atm}^{-1} \text{ m}^{-3}$  is required to offset the biological influence. The change in solubility with temperature ranges from  $-0.92$  to  $-0.6$  for SSTs of  $20\text{--}30^\circ\text{C}$ . Thus the increase in ingassing flux due to biological activity is equivalent to the increase in outgassing flux caused by an increase in SST of  $\sim 0.5\text{--}0.7 \text{ K}$ . Therefore if surface nutrient is available it is possible that biological activity could eliminate the temperature-induced increase in outgassing for  $\Delta\text{SST} \leq \sim 0.7 \text{ K}$ . However, since diurnal warming generally occurs when the ocean is strongly stratified, these are the times when there is less surface nutrient available and biological activity is probably much lower than we have estimated.

The analysis presented herein required the Taka02  $p\text{CO}_{2w}$  climatology to be interpolated over the shelf sea regions. This is not ideal but, in the absence of a shelf sea  $p\text{CO}_{2w}$  climatology, was the only approach. Figure 8 shows examples of the interpolation results for January and July. The method appears sensible through the Mediterranean Sea and around the coasts but very high values are estimated in the Red Sea (NE Africa) in July due to the high values in the Arabian sea in the Taka02 climatology. Whether or not the fluxes predicted over these regions are reasonable is unknown, however, since our concern is the difference in flux caused by SST variability, this is not critical. Similarly the Taka02 climatology is referenced to the year 1995, therefore the fluxes shown in this study are computed using driving data from 2005–2006 but can not be thought to be the actual fluxes for this period as the  $p\text{CO}_2$  fields have undoubtedly changed since 1995.

Finally, there is the issue of missing data. Since satellites measurements of SST are not possible through cloud, there are many missing data points. In fact in some

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regions there may not be a single satellite measurement for a month (see Fig. 7). Therefore, although the extreme diurnal warming events will occur under clear-sky conditions there may be more moderate warming events that are missed due to cloud.

## 6 Conclusions

5 Diurnal variations in SST have a significant impact on CO<sub>2</sub> flux over the SEVIRI disk region (central Atlantic ocean and Mediterranean). Including diurnal variability in SST increases the mass net flux out of the ocean from 9.6 Tg C a<sup>-1</sup> to 30.4 Tg C a<sup>-1</sup>. At a local scale, average monthly fluxes may be increased by up to ~0.2 mol CO<sub>2</sub> m<sup>-2</sup> a<sup>-1</sup>. This is due to the decrease in solubility associated with ΔSST rather than covaria-  
10 tions between diurnally varying variables (which causes a decrease in outgassing of ~1 Tg C a<sup>-1</sup>). Therefore, it is important that the additional outgassing of CO<sub>2</sub> due to ΔSST is accounted for, but it may be estimated using monthly-averaged values of ΔSST along with a foundation SST rather than high resolution SST data.

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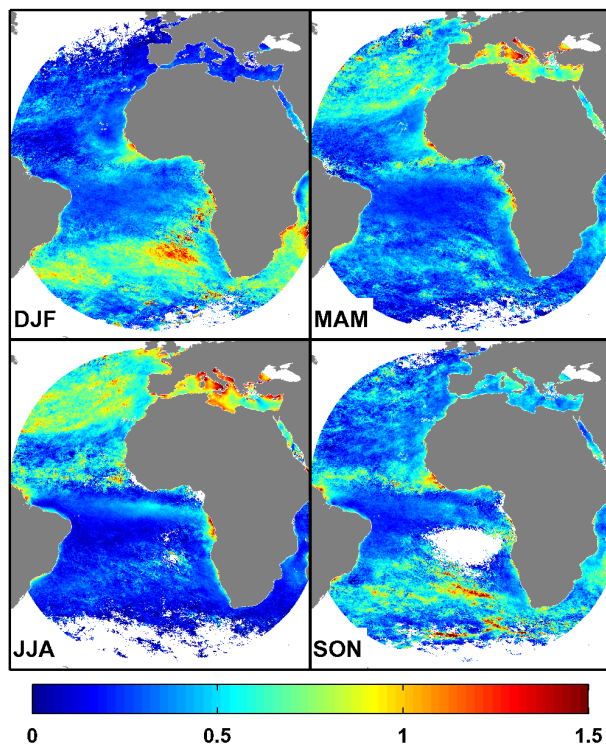
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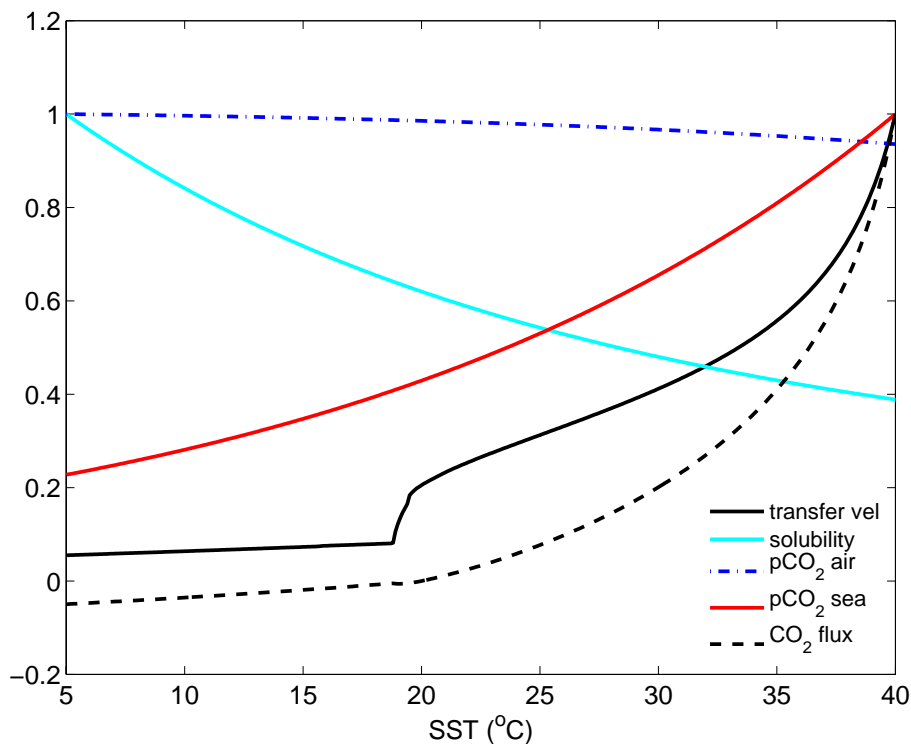


**Fig. 1.** Seasonal means of the mean daily peak  $\Delta$ SST (K) calculated from SEVIRI observations from June 2004 to May 2007, for northern winter (DJF), spring (MAM), summer (JJA) and autumn (SON). Regions with no valid data are marked white.

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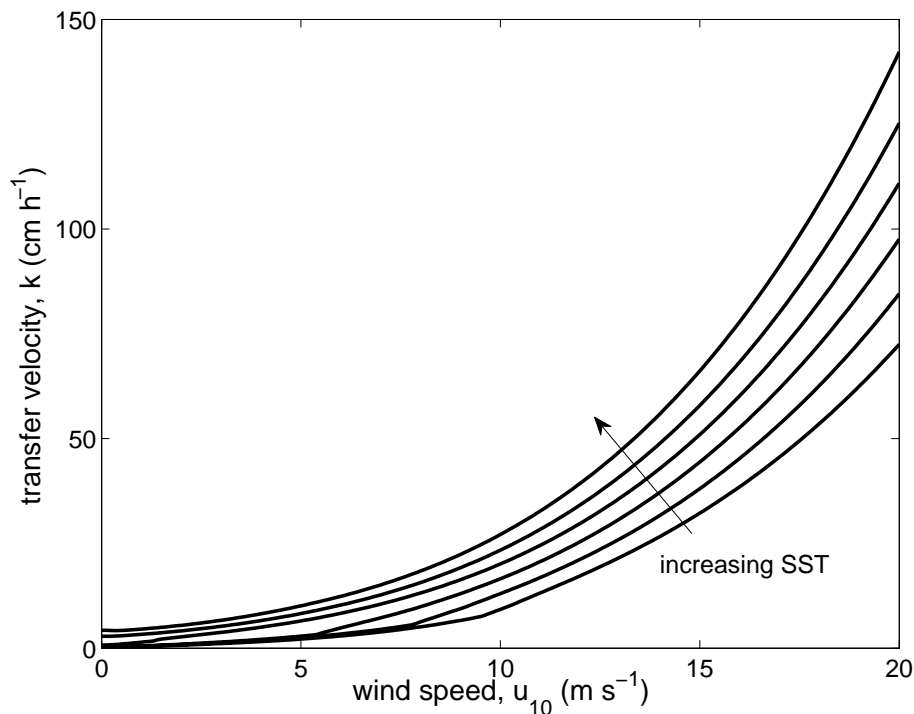


**Fig. 2.** Normalised (by maximum value) variables as a function of SST, with  $u_{10}=2 \text{ m s}^{-1}$ ,  $P=1000 \text{ mb}$ ,  $SSI=1000 \text{ W m}^{-2}$ ,  $DLI=300 \text{ W m}^{-2}$ ,  $\text{salinity}=35$ ,  $T_{\text{air}}=20^\circ\text{C}$ ,  $MLD=20 \text{ m}$ ,  $T_{\text{dew}}=15^\circ\text{C}$ ,  $[\text{CO}_2]_{\text{air}}=0.35 \mu \text{ atm mb}^{-1}$ , and the reference  $p\text{CO}_{2w}=350 \mu \text{ atm}$  at  $20^\circ\text{C}$ .

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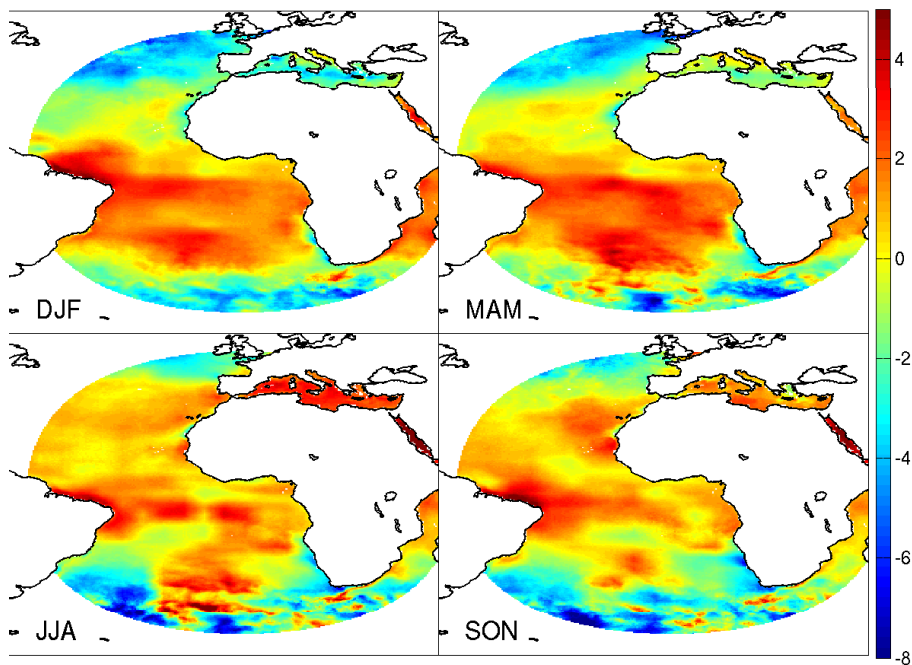
**Fig. 3.** Transfer velocity as a function of wind speed for SST=5, 10, 15, 20, 25, 30°C, with  $P=1000$  mb, SSI=1000 W m<sup>-2</sup>, DLI=300 W m<sup>-2</sup>, salinity=35,  $T_{\text{air}}=20^{\circ}\text{C}$ , MLD=20 m,  $T_{\text{dew}}=15^{\circ}\text{C}$ .

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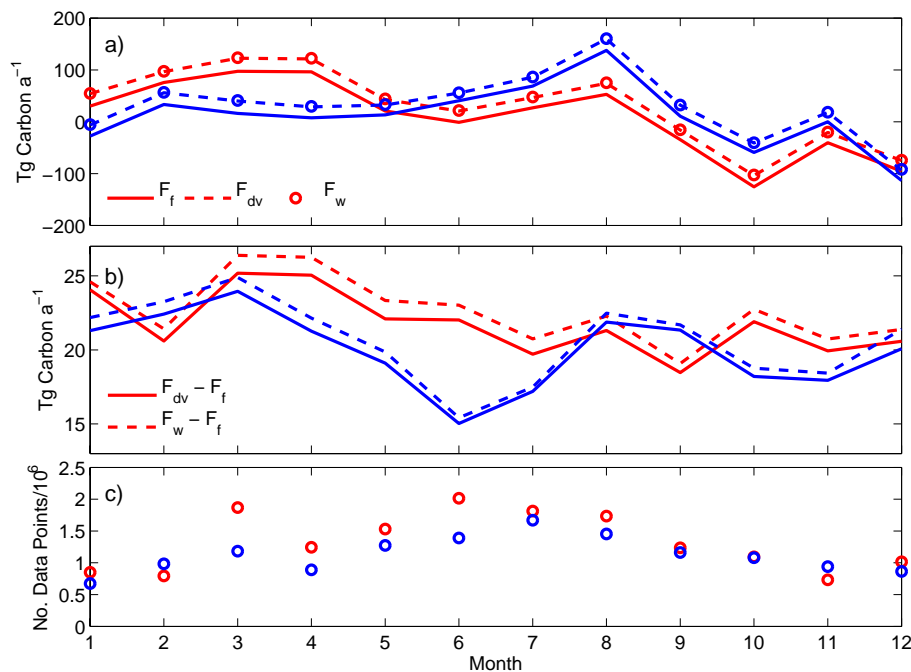


**Fig. 4.** Average net CO<sub>2</sub> flux (mol CO<sub>2</sub> m<sup>-2</sup> yr<sup>-1</sup>) for northern winter (DJF), spring (MAM), summer (JJA) and autumn (SON) during 2005-2006 using foundation SST ( $T_f$ ). Positive flux indicates outgassing of CO<sub>2</sub> from the ocean.

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**Fig. 5.** (a) Monthly-averaged mass flux of carbon over SEVIRI disk computed using the foundation SST ( $F_f$ ; solid line), including diurnal covariations ( $F_{dv}$ ; dashed line) and including monthly-averaged diurnal warming ( $F_w$ ; circles); (b) Difference in mass flux caused by including diurnal variations (solid line) and monthly-averaged warming (dashed line); (c) Number of satellite-derived  $\Delta\text{SST}$  data points available each month. 2005 is shown in red and 2006 in blue.

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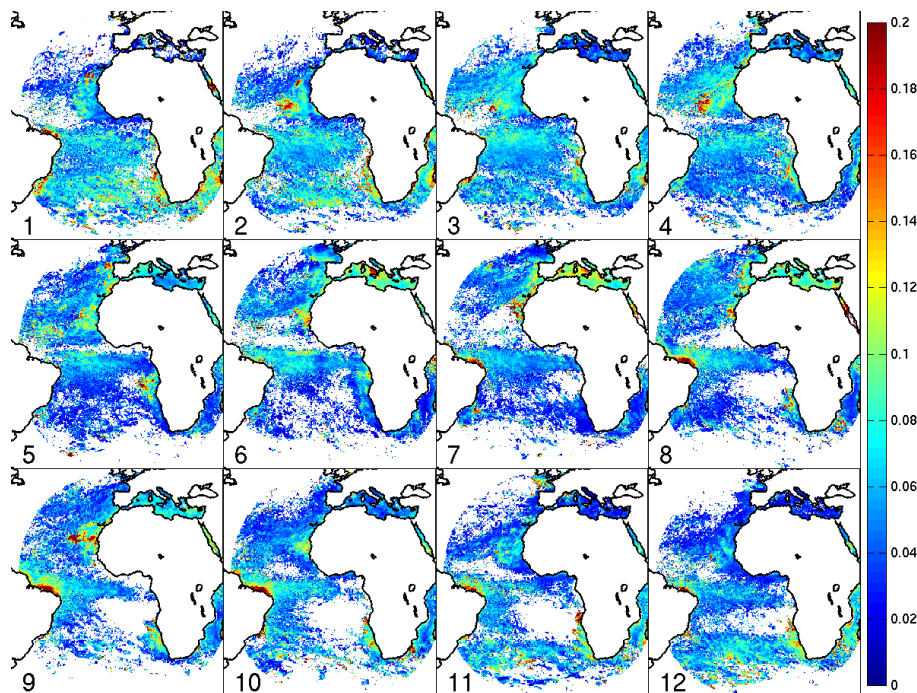
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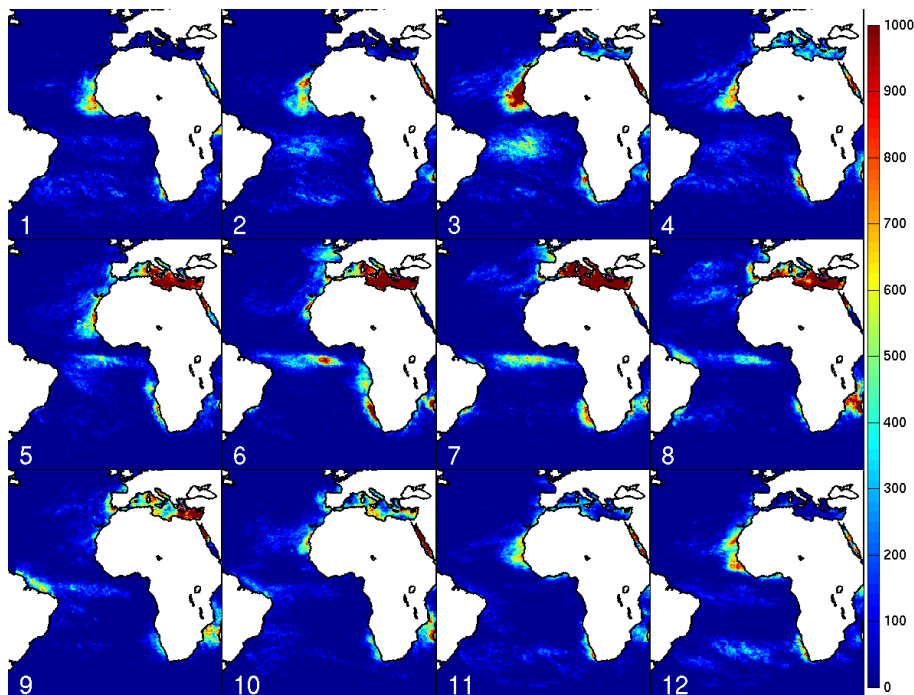
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**Fig. 6.** Average difference in the net CO<sub>2</sub> flux ( $\text{mol CO}_2 \text{ m}^{-2} \text{ yr}^{-1}$ ) averaged over each month (numbered) over 2005–2006 caused by diurnal variations in SST ( $F_{dv} - F_f$ ). Positive values indicate an increase in outgassing of CO<sub>2</sub> from the ocean. White regions indicate missing data.

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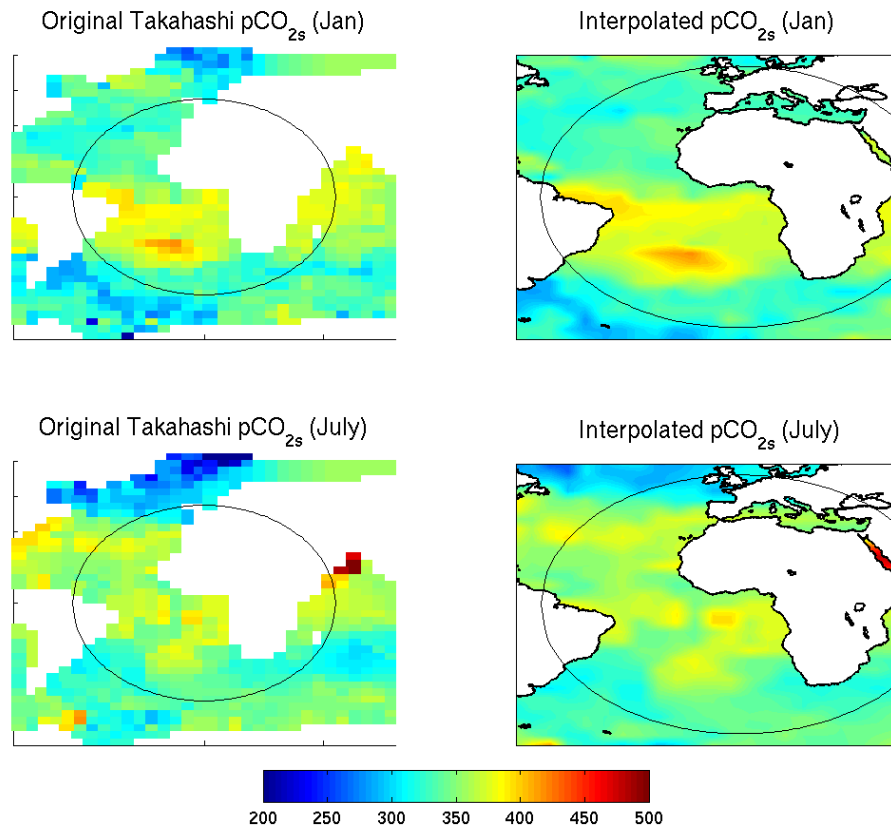


**Fig. 7.** Number of data points available to compute fluxes. Colorscale is saturated at 1000 to show data distribution by in data rich areas numbers go up to around 2600.

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**Fig. 8.** Interpolating  $p\text{CO}_{2w}$  from Taka02 for the reference year 1995 over the shelf seas for January (top row) and July (bottom row). Black circle indicates circumference of SEVIRI disk.

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